

## Synthesis and characterization of magneto electric coupling in nanostructured BiFeO<sub>3</sub> and BiFe<sub>1-x</sub>M<sub>x</sub>O<sub>3</sub> (M = Mn, Co, etc.) thin films for spintronics applications

P. Sateesh <sup>a, \*</sup>, A. Rajesh <sup>a</sup>, N. Maramu <sup>b</sup>, N. Hemanth <sup>a</sup>, K. Gopinath <sup>c</sup>,  
P. Sivakumar <sup>d</sup>, T. Satya <sup>e</sup>, G. K. Sivasankara Yadav <sup>f</sup>, T. Ramprasad <sup>a</sup>, M. Rajesh <sup>a</sup>

<sup>a</sup> Department of Physics, St. Peter's Engineering College, Secunderabad-500043

<sup>b</sup> Department of Physical science, Kakatiya Institute of technology and science, Warangal-506015

<sup>c</sup> Department of chemistry, St. Peter's Engineering College, Secunderabad-500043

<sup>d</sup> Department of Physics, Kaveri University, Gowraram- 507211, India.

<sup>e</sup> Department of H&BS, G. Pulla Reddy Engineering College (Autonomous), Kurnool-518007

<sup>f</sup> Department of Physics, Rayalaseema university, Kurnool-518007

In a recent study, researchers employed sol-gel deposition techniques to fabricate both pure BiFeO<sub>3</sub> and heterostructured BiFeO<sub>3</sub>/BiFe<sub>1-x</sub>M<sub>x</sub>O<sub>3</sub> thin films. The aim was to enhance the magnetic properties of BiFeO<sub>3</sub> through doping with Co<sup>3+</sup> and Mn<sup>3+</sup>. The team successfully prepared pure BiFeO<sub>3</sub> (BFO) and doped BiFe<sub>1-x</sub>M<sub>x</sub>O<sub>3</sub> films with varying concentrations of Mn and Co ions. XRD and Raman spectroscopy studies demonstrated that Co<sup>3+</sup> and Mn<sup>3+</sup> ions successfully substituted Fe<sup>3+</sup> ions within the rhombohedral lattice structure of BFO. The pure BiFeO<sub>3</sub> films showed stable and consistent responses when exposed to light, achieving an open circuit potential of -0.18 V and a short circuit photocurrent density of -0.024 mA/cm<sup>2</sup>. Additionally, the incorporation of Co<sup>3+</sup> and Mn<sup>3+</sup> ions led to a decrease in the average particle size. Recent research using Transmission Electron Microscopy (TEM) data has demonstrated that replacing Co<sup>3+</sup> and Mn<sup>3+</sup> with Fe<sup>3+</sup> in Perovskite structures greatly diminishes oxygen defects. This substitution also raises the concentration of Fe<sup>3+</sup> ions at the B-site. Complementary X-ray Photoelectron Spectroscopy (XPS) analysis confirms these results, showing improved magnetic properties in the thin film. Notably, the BFMn<sub>0.08</sub>O thin film exhibits a significant increase in both saturation and remanent magnetization at room temperature. Additionally, computational calculations reveal that CO exhibits a strong preference for adsorption on Mn-doped BiFeO<sub>3</sub>, with an adsorption energy calculated at -1.0494 eV.

(Received May 7, 2025; Accepted August 21, 2025)

**Keywords:** BiFeO<sub>3</sub> thin film, Perovskite grid, Photo illumination, Spintronics

### 1. Introduction

Multiferroic materials have become a focal point of research due to their unique physical characteristics and their potential use in devices such as sensors, memory storage, and spintronic components. These materials are categorized into two primary groups: type-I and type-II multiferroics. Type-I multiferroics are notable for their strong ferroelectric properties, though the temperatures at which their magnetic and ferroelectric transitions occur can often surpass room temperature. In these materials, the magnetoelectric (ME) coupling tends to be weak, as it involves interactions between magnetization and polarization. Ferroelectricity in these compounds can arise through mechanisms such as charge ordering or the displacement of ions or ion groups, commonly observed in transition metal oxides like hexagonal RMnO<sub>3</sub>. Type-II multiferroic materials, also known as "spin driven ferroelectrics," exhibit magnetic ordering that directly induces ferroelectricity by disrupting inversion symmetry. This phenomenon is attributed to their non-

---

\* Corresponding author:

<https://doi.org/10.15251/DJNB.2025.203.975>

collinear spin spiral structure, which generates the necessary symmetry breaking. In these materials, macroscopic electric polarization arises from the non-centrosymmetric long-range magnetic order. A prime example is orthorhombic  $\text{RMnO}_3$ , where the temperatures for ferroelectric and magnetic phase transitions are closely aligned, resulting in strong magneto electric (ME) coupling. These materials demonstrate both spontaneous polarization and antiferromagnetic ordering, with an antiferromagnetic Néel temperature around 643K and a high ferroelectric Curie temperature near 1103K. Research has explored enhancing the magnetic properties of  $\text{BiFeO}_3$  through B-site (Fe-site) cation substitution, thereby improving its multiferroic characteristics.

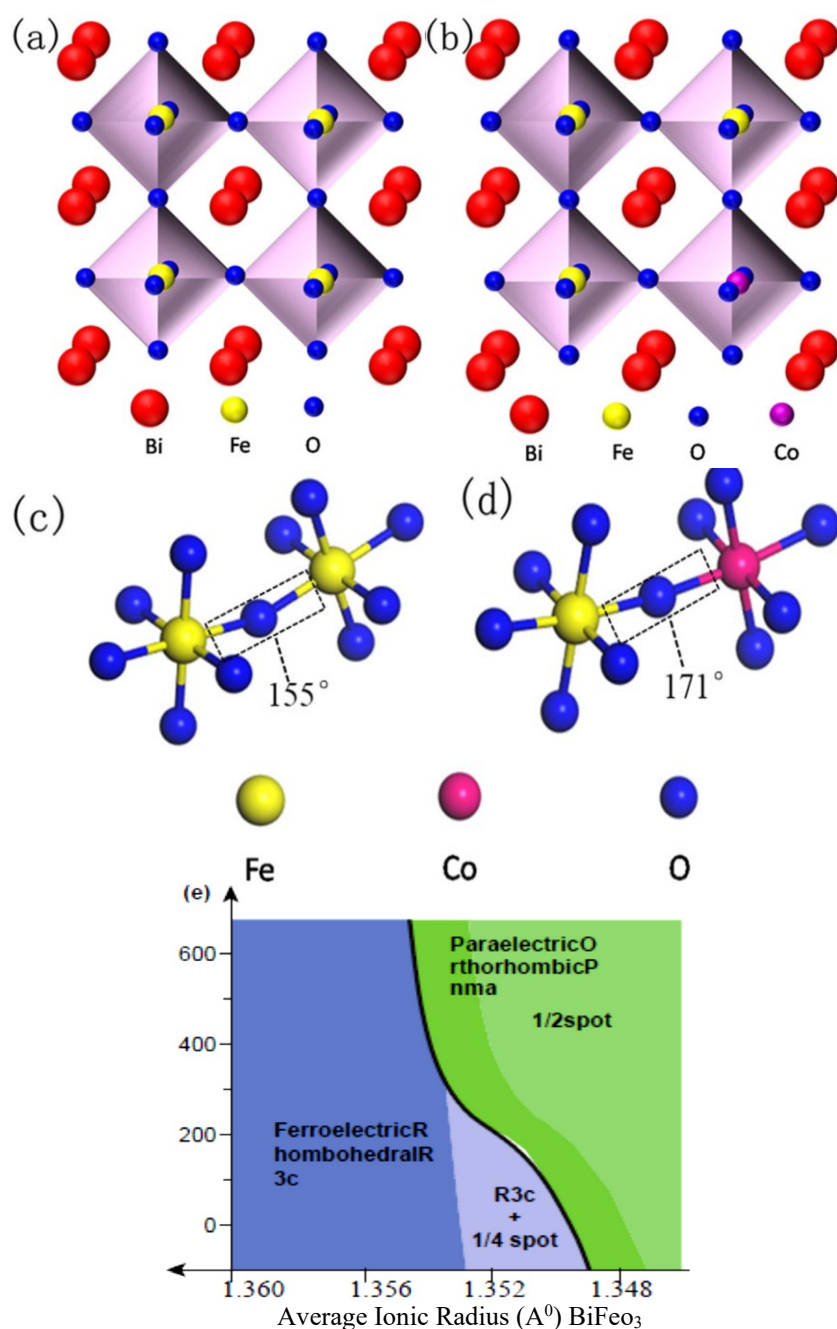


Fig. 1. The simulations included the Polycrystalline  $\text{BiFeO}_3$  (BFO) crystal structure and the Mn-Co switched BFO. Additionally, the study examined the super-swapping interactions between  $\text{Fe-O-Fe}$  and  $\text{Fe-O-Co}$ .

## 2. Experimental details

Thin films of pure BiFeO<sub>3</sub> (BFO) and doped BiFe<sub>1-x</sub>M<sub>x</sub>O<sub>3</sub> (where M represents elements like Mn or Co, and x values are 0.01, 0.03, 0.06, 0.08, and 0.10) were synthesized on silicon substrates utilizing the sol-gel method. The starting materials, including bismuth nitrate [Bi(NO<sub>3</sub>)<sub>3</sub>•5H<sub>2</sub>O], ferric nitrate [Fe(NO<sub>3</sub>)<sub>3</sub>•9H<sub>2</sub>O], and cobalt (II) nitrate hexahydrate [Co(NO<sub>3</sub>)<sub>2</sub>•6H<sub>2</sub>O], were sourced from Hyderabad, India, and manganese acetate tetrahydrate [Mn(C<sub>2</sub>H<sub>3</sub>O<sub>2</sub>)<sub>2</sub>•4H<sub>2</sub>O]. These were fully dissolved in a mixture of glacial acetic acid and ethylene glycol methyl ether to form the precursor solution. Following the precise addition and vigorous agitation of C<sub>6</sub>H<sub>8</sub>O<sub>7</sub>, C<sub>2</sub>H<sub>6</sub>O<sub>2</sub>, and C<sub>2</sub>H<sub>7</sub>NO, the solution reached clarity and matured to a 0.4 mol/L concentration over 28 hours. The spin-coating technique involved two phases: an initial spin at 800 rpm for 3 seconds followed by 5000 rpm for 18 seconds. The wet films underwent a pre-annealing process at a temperature of 390°C for a duration of 8 minutes. This procedure was carried out for each layer, allowing for the construction of a film with moderate thickness. The samples underwent a final crystallization step in air at 600°C for one hour before cooling to ambient temperature, resulting in optimal thin films. Figure 2 illustrates the experimental process.

Figure 2 illustrates the procedural diagram for synthesizing BiFe<sub>1-x</sub>M<sub>x</sub>O<sub>3</sub> thin films, where M represents elements such as Mn or Co. The diagram covers the doping concentrations of x = 0.01, 0.03, 0.06, 0.08, and 0.10, highlighting the integration process for these materials.

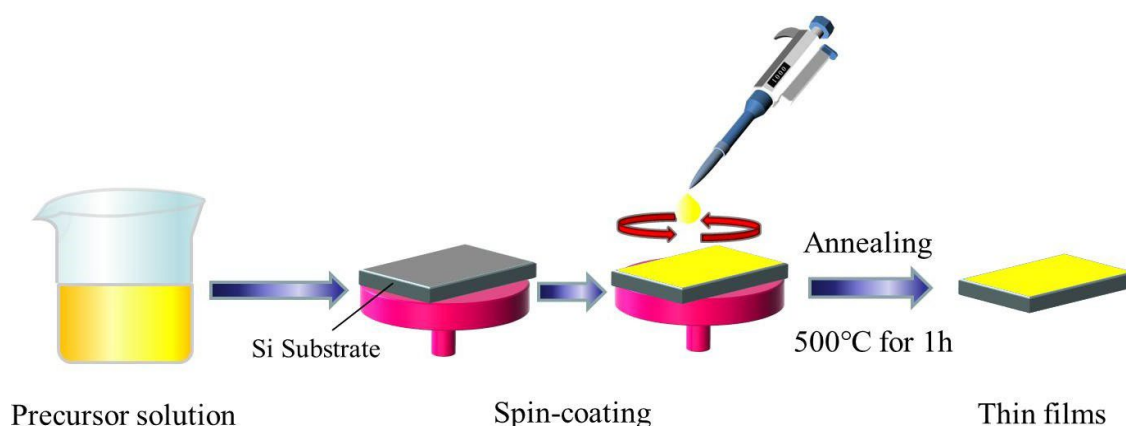


Fig. 2. (e) Phase diagram for rare earth replaced BFO thin films that has been proposed.

## 3. Results and discussion

To aid in understanding doping principles and super-swapping interactions, figures 3 and 4 simulate the crystal structures of both doped BiFe<sub>1-x</sub>M<sub>x</sub>O<sub>3</sub> and pure BFO. The BFO composite thin film, which belongs to the space group R3C, displays a deformed rhombohedral Perovskite structure. This structure is characterized by iron ions at the center of an FeO<sub>6</sub> octahedron surrounded by six oxygen ions, while bismuth ions form a cube encasing the octahedron. In particular, in Co<sup>3+</sup>/Mn<sup>3+</sup>-doped BFO films, one oxygen anion connects two adjacent FeO<sub>6</sub>/CoO<sub>6</sub>/MnO<sub>3</sub> octahedrons. The difference in size between the B-site cations, specifically Fe<sup>3+</sup> and Co<sup>3+</sup>/Mn<sup>3+</sup>, causes both dimensional and chemical strain reactions. These reactions disrupt the modulation of the spin cycloid spiral. As a result, some Fe<sup>3+</sup> ions are randomly substituted at the B-site. The substitution of Co and Mn ions for Fe ions is clearly demonstrated in Figures 3 and 4 [21]. This substitution leads to the formation of high-energy Co/Mn–O bonds, enhancing the inclination angle of the super-exchange interaction. As a result, the improved super-exchange interaction of Fe–O–Fe (or Mn/Co) contributes to the enhanced magnetic properties of BiFe<sub>1-x</sub>M<sub>x</sub>O<sub>3</sub>, as detailed in the following sections.

The X-ray diffraction (XRD) patterns of the  $\text{BiFe}_{1-x}\text{MxO}_3$  composite thin films, as shown in Figures 3 and 4, exhibit diffraction peaks that align with the standard XRD profile for the rhombohedral phase (space group  $R3c$ ), as referenced by JCPDS71-2494 and 53-0767. An impurity peak at  $28.8^\circ$ , attributed to  $\text{Bi}_2\text{Fe}_4\text{O}_9$ , is also observed [17,21]. Notably, the disappearance of the (021) peak and shifts in the (012) and (110) peaks around  $2\theta = 56^\circ$ – $57.1^\circ$  indicate a structural transition in the BFO films, from rhombohedral to tetragonal, influenced by Mn doping, consistent with findings from tetragonal  $\text{BiMnO}_3$  (PDF#53-0767) and corroborated by earlier studies [18].

According to XRD analysis, the intrinsic structures and properties of  $\text{BiFe}_{1-x}\text{MxO}_3$  are unlikely to be significantly affected by the presence of  $\text{Bi}_2\text{Fe}_4\text{O}_9$ , owing to its paramagnetic nature and low concentration in the samples [17,21]. The magnified XRD patterns of Co and Mn-substituted thin films at different concentrations are shown in Figure 3b, focusing on the  $2\theta$  range of  $32^\circ$ . An increase in  $\text{Co}^{3+}/\text{Mn}^{3+}$  concentration results in the diffraction peaks for (104) and (110) moving to higher angles. This shift indicates the effective integration of Co ions into the BFO host lattice. This shift is attributed to the smaller ionic radius of  $\text{Co}^{3+}/\text{Mn}^{3+}$  compared to  $\text{Fe}^{3+}$ , resulting in structural deformation and lattice contraction.

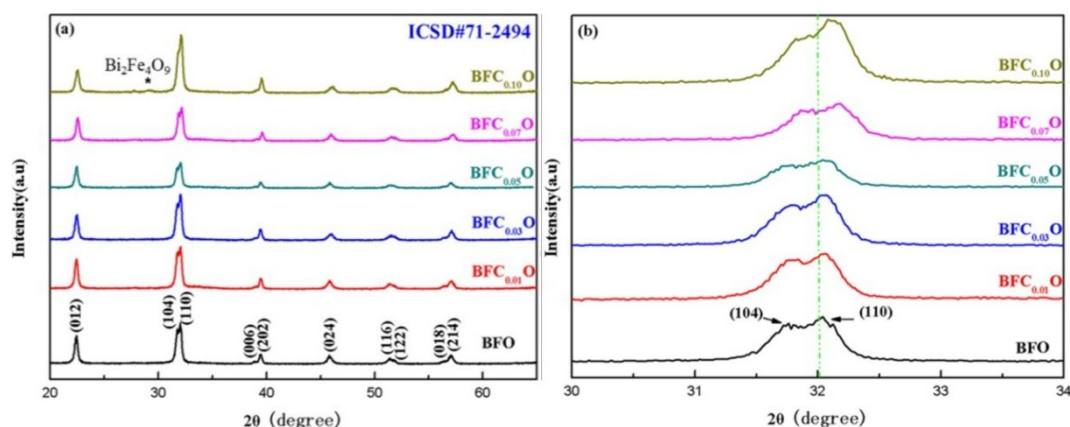


Fig. 3. XRD Patterns of  $\text{BiFe}_{1-x}\text{MxO}_3$  ( $M = \text{Co}$ ) and Pure BFO thin films.

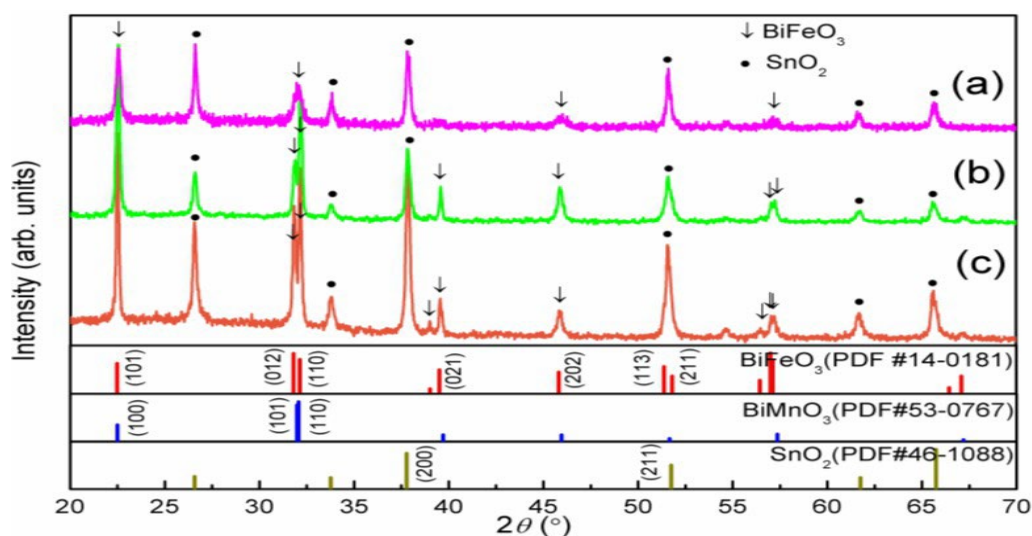
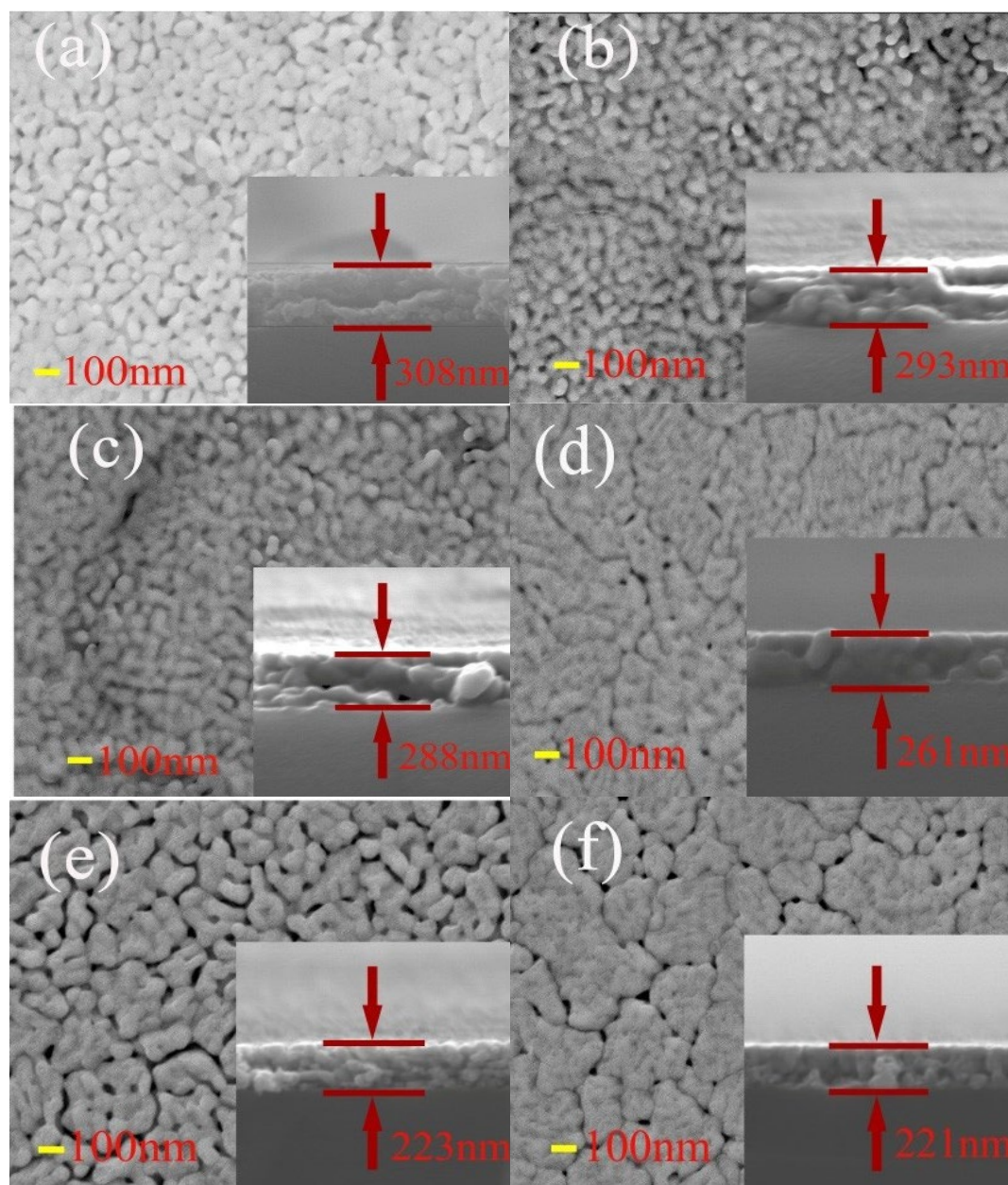


Fig. 4. XRD Patterns of  $\text{BiFe}_{1-x}\text{MxO}_3$  ( $M = \text{Mn}$ ) thin films.



Figure 5a presents the surface and cross-sectional SEM pictures of the  $\text{BiFe}_{1-x}\text{MxO}_3$  ( $\text{M}=\text{Mn}, \text{Co}, \text{etc.}$ ) thin films are shown. The Co-substituted thin films are demonstrated to have a consistent surface shape, observable aggregation, and good adhesion to the silicon substrates without apparent separation. Additionally,  $\text{BiFe}_{1-x}\text{MxO}_3$  thin film thicknesses are 308, 293, 288, 261, 223 and 221 nm, respectively, where M is Mn, Co, etc. It appears that Co/Mn doping results in a uniformly dense and grainy surface because the thickness obviously reduces as the amount of Co doping increases [26, 27].



*Fig. 5. SEM images illustrating the surface and cross-sectional morphologies of various thin films. These include (a) BFO, (b) BFC/Mn0.010, (c) BFC/Mn0.030, (d) BFC/Mn0.060, (e) BFC/Mn0.080, and (f) BFC/Mn0.100. These images provide valuable insights into the structural characteristics of each film composition.*

#### 4. Magnetic properties

In a magnetic field ranging from 0 to 15 k, we examined how Co and Mn doping affects the magnetic properties of BFO and BiFe 1-xMxO<sub>3</sub> thin films (M = Mn, Co, etc.). Figure 8a illustrates the saturated magnetic hysteresis loops observed in both BFO and Co-Mn doped thin films'-Mn tampering with concentration increased saturation magnetization ( $M_s$ ) as well as remanent magnetization ( $M_r$ ). The materials BFO, BFC/Mn 0.01 O, BFC/Mn 0.03 O, BFC/Mn 0.05 O, BFC/Mn 0.07 O, and BFC/Mn 0.10 O exhibit the following magnetization values: 21.38 emu/cm<sup>3</sup> and 2.06 emu/cm<sup>3</sup>, 25.57 emu/cm<sup>3</sup> and 2.40 emu/cm<sup>3</sup>, 34.17 emu/cm<sup>3</sup> and 2.60 emu/cm<sup>3</sup>, 34.48 emu/cm<sup>3</sup> and 3.00 emu/cm<sup>3</sup>, 42.38 emu/cm<sup>3</sup> and 3.68 emu/cm<sup>3</sup>, and 47.84 emu/cm<sup>3</sup> and 3.43 emu/cm<sup>3</sup>, respectively. Figure 9 demonstrates the variations in Co doping concentrations between  $M_s$  and  $M_r$ . In the case of  $M_s$ , the values increase in a linear fashion, whereas  $M_r$  initially increases but then decreases as the levels of Co<sup>3+</sup> and Mn<sup>3+</sup> surpass  $x = 0.12$ . The highest recorded value of  $M_r$  is 4.64 emu/cm<sup>3</sup> at a cobalt concentration of 0.08. In contrast, the maximum value of  $M_s$  is 47.88 emu/cm<sup>3</sup>, which occurs when the doping concentration reaches 0.10.

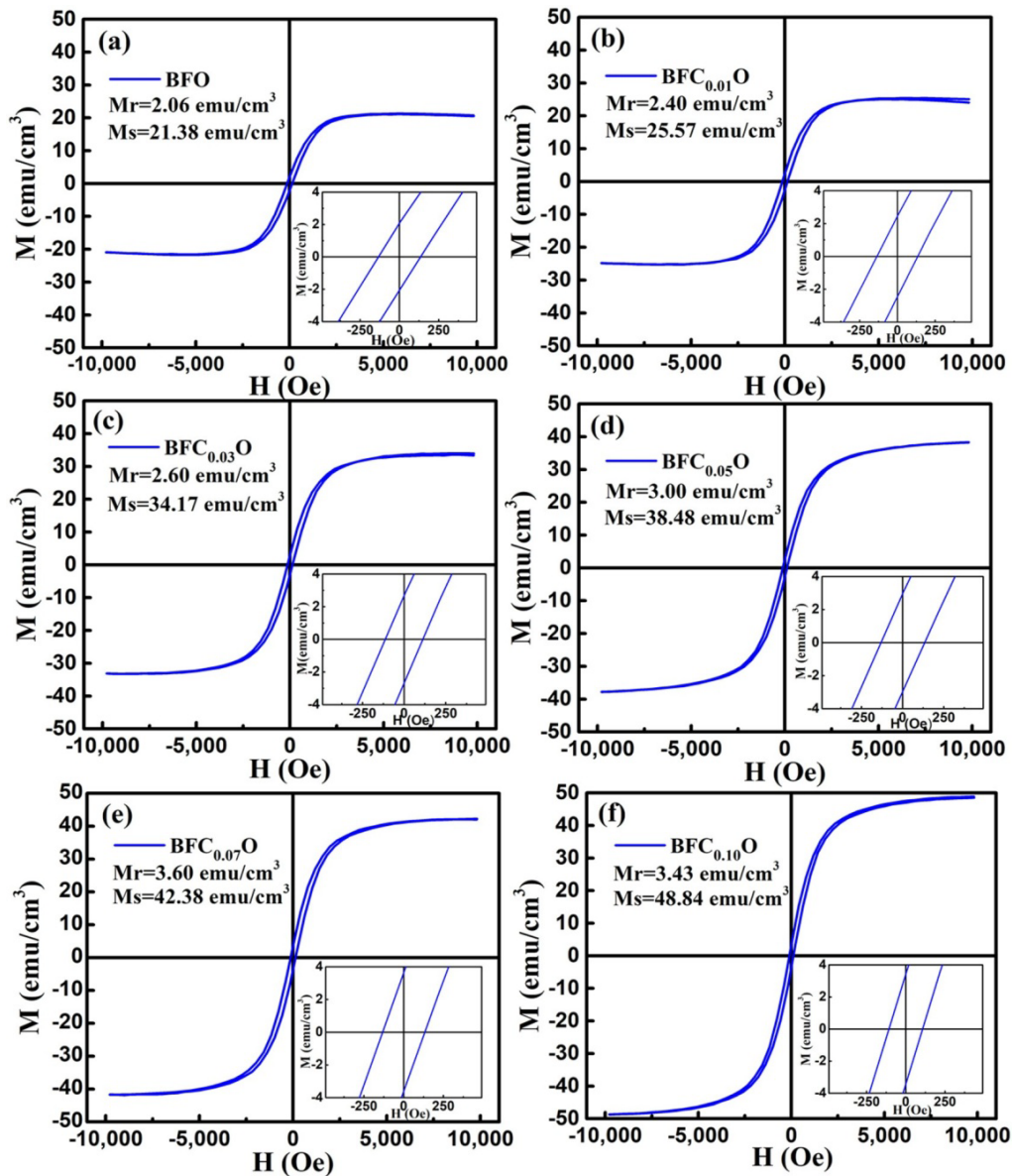


Fig. 8. Presents the magnetic hysteresis ( $M$ - $H$ ) loops for various room temperature thin films. These include (a) BFO, (b) BFC/Mn0.01O, (c) BFC/Mn0.03O, (d) BFC/Mn0.05O, (e) BFC/Mn0.08O, and (f) BFC/Mn0.10O. Each loop provides insight into the magnetic properties of the films as influenced by differing manganese oxide (MnO) concentrations.

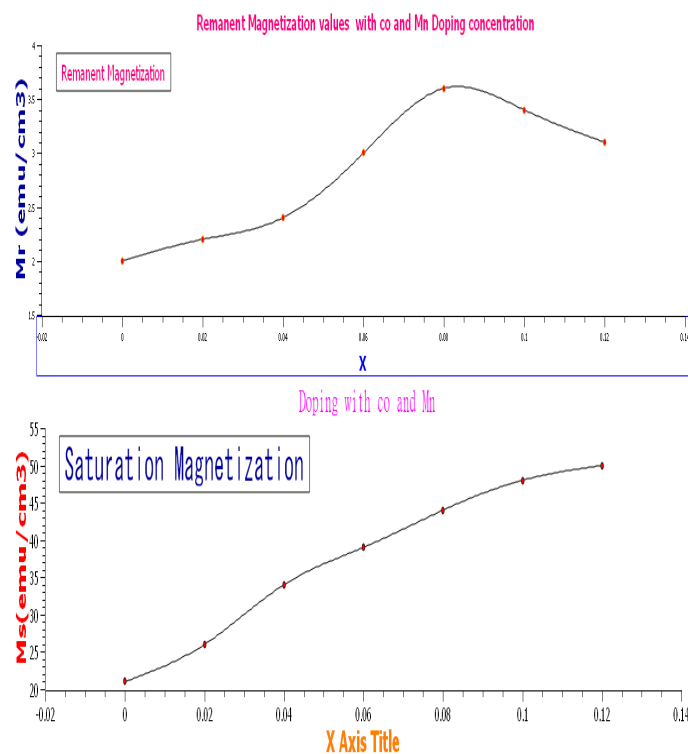


Fig. 9. Variations in saturation magnetization ( $M_s$ ) and remanent magnetization ( $M_r$ ) values with concentrations of Co and Mn stupefy.

#### 4.1. Optical properties

The optical characteristics of BFO nanoparticles Current density for BiFe 1 -xMxO3 and BFO thin films (M = Mn, Co, etc., x = 0.01, 0.03, 0.06, 0.08, and 0.10) (b) The photocurrent in BiFe 1-xMxO3 thin-film heating and cooling states at varying light intensities. (c) The output current with BFO and BiFe 1 -xMxO3 (M = Mn, Co) thin films as a function of temperature.

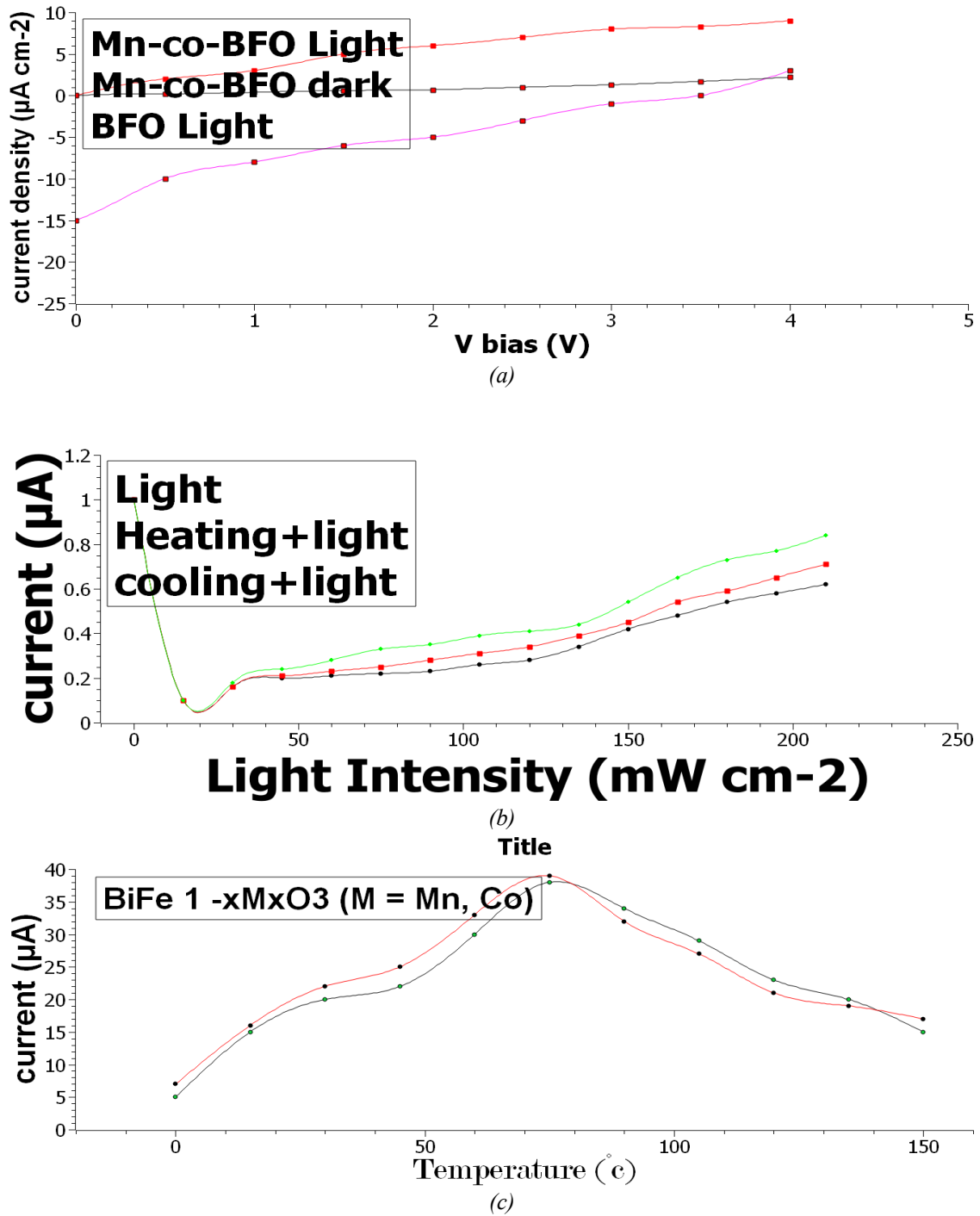


Fig. 10. (a) Photovoltaic performance for different current densities (b) Photovoltaic current performance for different light intensities (c) Photovoltaic current performance for different temperatures

#### 4.2. Photo catalysis

With a greater light, photocurrent was seen in the heating condition. It was discovered that the cooling state increased the photocurrent under weak light illumination. The photovoltaic properties of  $\text{BiFe } 1-x\text{MxO}_3$  ( $\text{M} = \text{Mn, Co, etc.}, x = 0.01, 0.03, 0.06, 0.08, \text{ and } 0.10$ ) thin film materials can be improved using a unique approach, as this work showed that temperature gradients can be used to adjust photovoltaic performance. It is possible for the photocurrent to increase and then drop as the temperature rises. The photo-sensing device's modulated energy bandgap and barrier height are the cause of its temperature-dependent performance.



The aforementioned studies offer fresh perspectives on the photovoltaic effect and offer recommendations for creating high-performance  $\text{BiFe}_{1-x}\text{MxO}_3$  ( $\text{M} = \text{Mn}, \text{Co}$ ) power devices.

#### **4.3. Applications of $\text{BiFe}_{1-x}\text{MxO}_3$ ( $\text{M} = \text{Mn}, \text{Co}$ , etc.):**

Recent advancements in BFO nanomaterials, specifically those doped with manganese (Mn) and cobalt (Co), have demonstrated promising capabilities in the realm of high-density ferroelectric devices. These materials, which exhibit both ferroelectricity and ferromagnetism at room temperature, offer a unique opportunity to leverage magnetoelectric coupling for advanced applications. Researchers have further explored this potential by developing a magneto resistive device through spin valve exchange, paving the way for innovative applications in electrically written spintronic devices. The study explored a device with a structure that demonstrated a 2% giant magnetoresistance signal, utilizing layers of STO/SRO/BFO-Mn-co/BFO/Mn-CoFeB/Cu/Mn-Co. Researchers found that applied voltages could effectively control the exchange bias. However, the weak thermal stability of the BFO/Mn-CoFeB interfaces was a significant barrier to enhancing the magnetoresistance signal. The analysis revealed that the main factor affecting thermal stability was the variation in oxygen vacancies, influenced by the BFO polarization direction. To achieve a higher magnetoresistance value of 4.4%, a new, efficient method was proposed.

#### **4.4. Application for spintronic and optic devices**

A) An enormous BFO-based magneto-resistive device. Source: b) swapping coupling with one the ferromagnetic over layer and ant ferromagnetic order for BFO/Mn-co-based spintronic devices, with permission. C) Unidirectional anisotropy using voltage control for Mn-CoFe/BFO heterostructures.

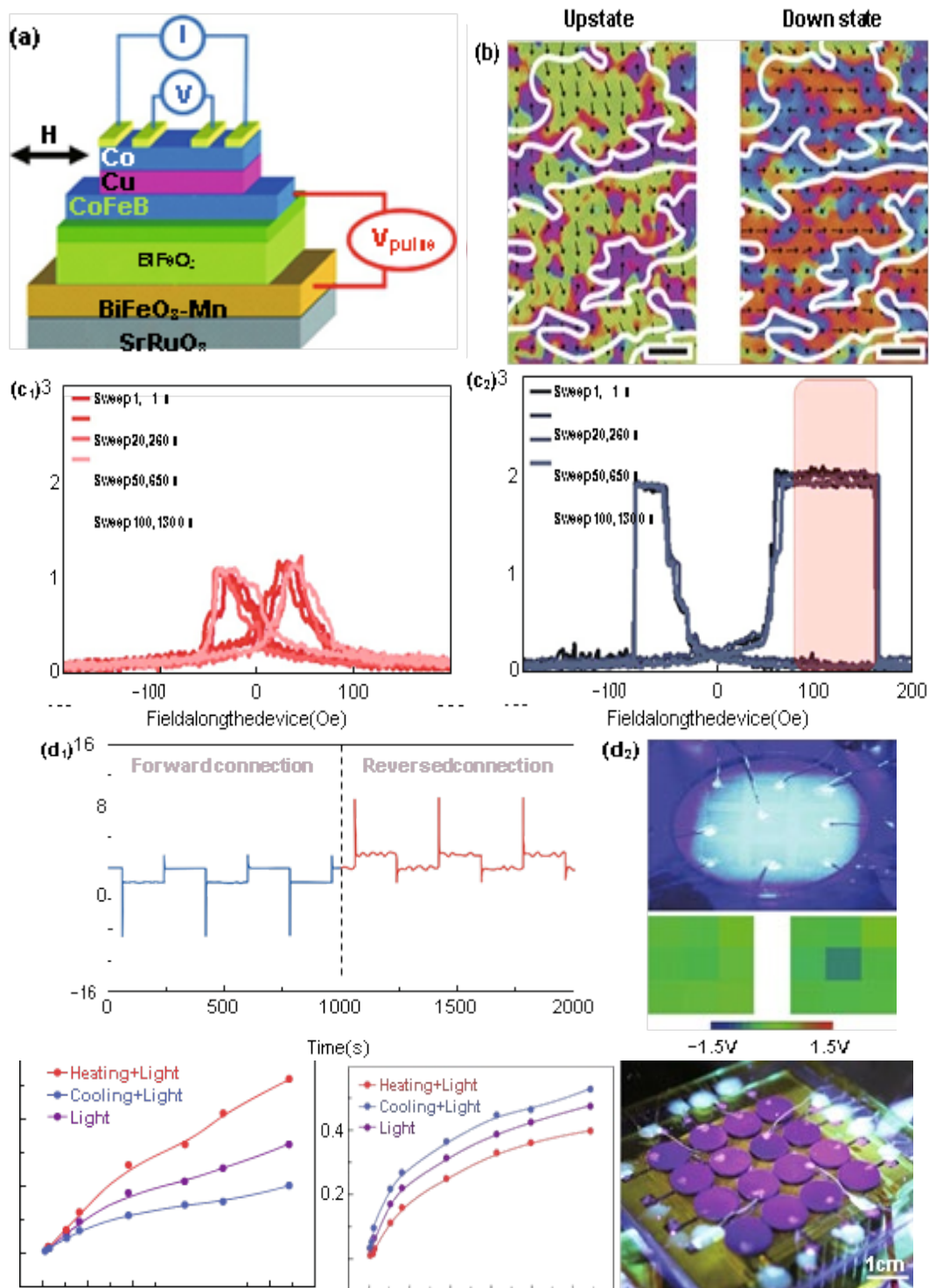


Fig. 11. Applications for spintronic and optic devices. Adapted by BFO

## 5. Conclusions

Recent advancements in nanomaterials have highlighted the exceptional properties of BFO, a multiferroic material known for its impressive characteristics such as high remnant polarization, magnetoelectric coupling at room temperature, and a relatively low bandgap.

In this study, Mn-Co modified BFO thin films were synthesized using a straightforward sol-gel technique. Comprehensive analyses were performed on the magnetic attributes, surface morphologies, and structural connections of the materials, revealing that all samples possess a rhombohedral polycrystalline structure. XRD and Raman spectroscopy confirmed the optimal incorporation of  $\text{Co}^{3+}$  and  $\text{Mn}^{3+}$  ions into the BFO lattice, while SEM imaging demonstrated that Mn-Co doping facilitates uniform particle growth and reduces grain size. Recent studies have shown that Co substitution in thin films can significantly enhance magnetic properties through noticeable agglomerations with smaller grain sizes. Vibrating Sample Magnetometer (VSM) evaluations reveal that these magnetic attributes are greatly improved, with the thin films of BF/Mn-Co0.08O and BF/Mn-Co0.10O achieving impressive  $M_r$  and  $M_s$  values of 3.68 and 47.84  $\text{emu/cm}^3$ , respectively.

This enhancement is attributed to factors such as reduced oxygen defects, denser accumulations, improved super exchange interactions, and the structure breakdown of spiral-spin systems. However, challenges remain in developing new devices as ferroelectric and magnetoelectric performances still require further optimization, despite the high polarizations observed in BF/Mn-Co.

## References

- [1] R. Guo, L. You, Y. Zhou, Z.S. Lim, X. Zou, L. Chen, R. Ramesh, J.L. Wang, *Nat. Commun.* 4, 1990 (2013); <https://doi.org/10.1038/ncomms2990>
- [2] J. Qi, N. Ma, Y. Yang, *Adv. Mater. Interfaces* 5, 1701189 (2018); <https://doi.org/10.1002/admi.201701189>
- [3] J. Qi, N. Ma, X.C. Ma, R. Adelung, Y. Yang, *ACS Appl. Mater. Interfaces* 10, 13712- 13719 (2018); <https://doi.org/10.1021/acsami.8b02543>
- [4] J. Ma, J. Ma, Q. Zhang, R. Peng, J. Wang et al., *Nat. Nanotechnol.* 13, 947-952 (2018); <https://doi.org/10.1038/s41565-018-0204-1>
- [5] F. Mushtaq, X. Chen, M. Hoop, H. Torlakcik, E. Pellicer et al., *iScience* 4, 236-246 (2018); <https://doi.org/10.1016/j.isci.2018.06.003>
- [6] H. Singh, K.L. Yadav, *Ceram. Int.* 41, 9285-9295 (2015); <https://doi.org/10.1016/j.ceramint.2015.03.212>
- [7] C.R. Foschini, M.A. Ramirez, S.R. Simoes, J.A. Varela, E. Longo, A.Z. Simoes, *Ceram. Int.* 39, 2185-2195 (2013); <https://doi.org/10.1016/j.ceramint.2012.08.083>
- [8] Dong, W.; Guo, Y.; Guo, B.; Liu, H.; Li, H.; Liu, H., *Mater. Lett.* 2012, 88, 140-142; <https://doi.org/10.1016/j.matlet.2012.08.006>
- [9] Peng, Z.; Liu, B., *Phys. Status Solidi A* 2014, 2, 1-4; <https://doi.org/10.1002/pssa.201431512>
- [10] Katiyar, R.K.; Sharma, Y.; Misra, P.; Puli, V.S.; Sahoo, S.; Kumar, A.; Scott, J.F.; Morell, G.; Weiner, B.R.; Katiyar, R.S., *Appl. Phys. Lett.* 2014, 105, 172904; <https://doi.org/10.1063/1.4900755>
- [11] Wang, Z.; Cao, D.; Wen, L.; Xu, R.; Obergfell, M.; Mi, Y.; Zhan, Z.; Nasori, N.; Demsar, J.; Lei, Y., *Nat. Commun.* 2016, 7, 10348; <https://doi.org/10.1038/ncomms10348>
- [12] Wu, J.; Fan, Z.; Xiao, D.; Zhu, J.; Wang, J., *Prog. Mater. Sci.* 2016, 84, 335-402; <https://doi.org/10.1016/j.pmatsci.2016.09.001>
- [13] Khomchenko, V.A.; Karpinsky, D.V.; Paixão, J.A., *J. Mater. Chem. C* 2017, 5, 3623-3629; <https://doi.org/10.1039/C7TC00833C>
- [14] Mao, W.; Wang, X.; Chu, L.; Zhu, Y.; Wang, Q.; Zhang, J.; Yang, J.; Li, X.; Wei, H. *Phys.*

- Chem. Chem. Phys. 2016, 18, 6399-6405; <https://doi.org/10.1039/C5CP07327H>
- [15] Dutta, D.P.; Tyagi, A., Appl. Surf. Sci. 2018, 450, 429-440; <https://doi.org/10.1016/j.apsusc.2018.04.202>
- [16] Qiu, H.; Chen, G.; Fan, R.; Cheng, C.; Hao, S.; Chen, D.; Yang, C., Chem. Commun. 2011, 47, 9648; <https://doi.org/10.1039/c1cc13707g>
- [17] Rong, Q.-Y.; Wang, L.-L.; Xiao, W.-Z.; Xu, L., Phys. B Condens. Matter 2015, 457, 1-4; <https://doi.org/10.1016/j.physb.2014.08.028>
- [18] Marzouki, A.; Harzali, H.; Loyau, V.; Gemeiner, P.; Zehani, K.; Dkhil, B.; Bessais, L.; Megriche, A., Mater. 2018, 145, 316-321; <https://doi.org/10.1016/j.actamat.2017.11.055>
- [19] Liu, Y.; Qi, J.; Zhang, Y.; Wang, Y.; Feng, M.; Zhang, J.; Wei, M.; Yang, J., Appl. Surf. Sci. 2018, 427, 745-752; <https://doi.org/10.1016/j.apsusc.2017.08.083>
- [20] Sahni, M.; Mukaherjee, S.; Hamid, A.; Kumar, D.; Chauhan, S.; Kumar, N. J. Mater. Sci. Mater. Electron. 2020, 31, 7798-7810; <https://doi.org/10.1007/s10854-020-03318-6>
- [21] Li, L.; Salvador, P.A.; Rohrer, G.S., Nanoscale 2014, 6, 24-42; <https://doi.org/10.1039/C3NR03998F>